

**ATMOSPHERIC CORRECTION  
of OCEAN IMAGERY**

**FOR**

**SeaWiFS and MODIS**

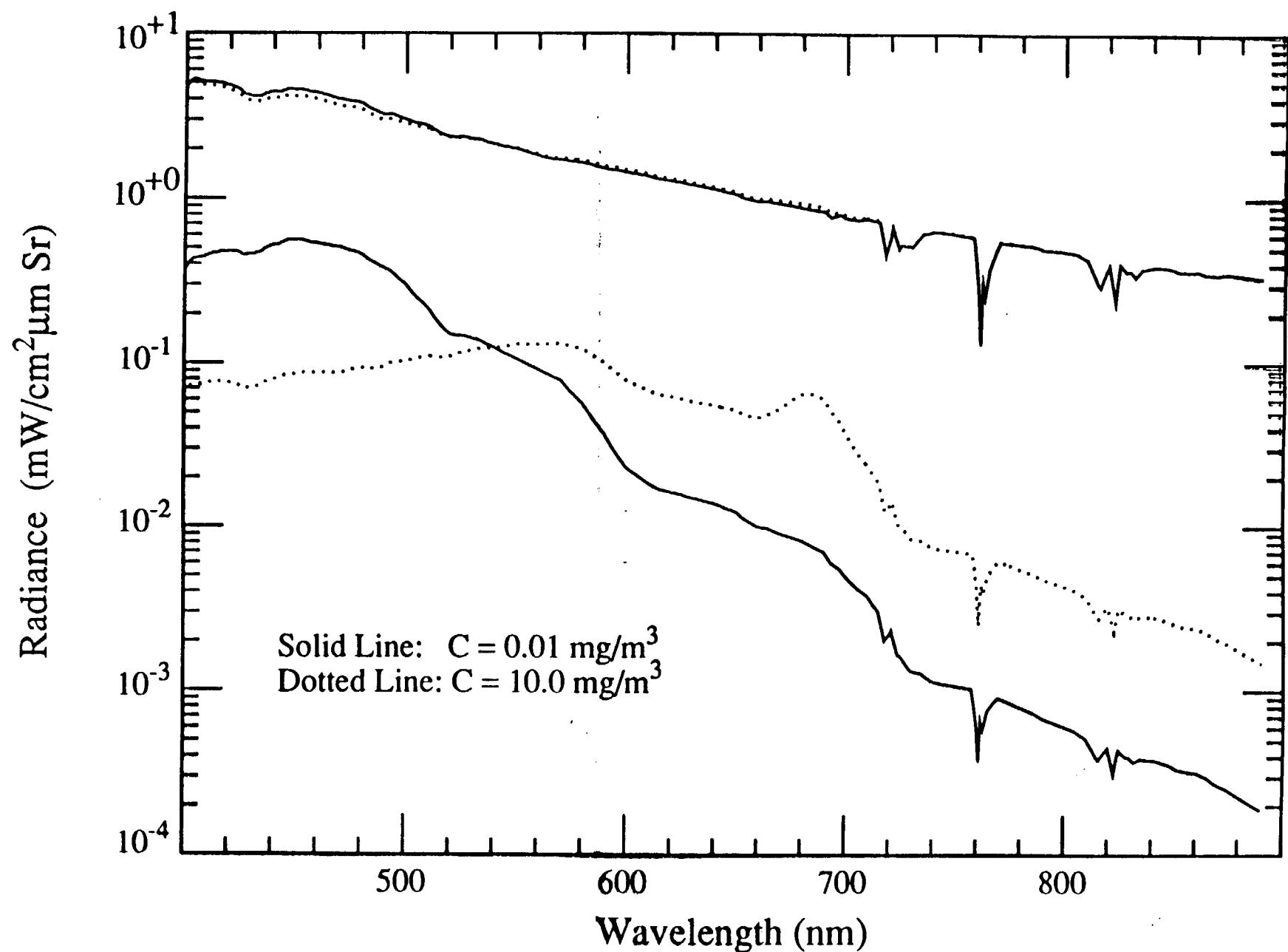
**Howard R. Gordon**

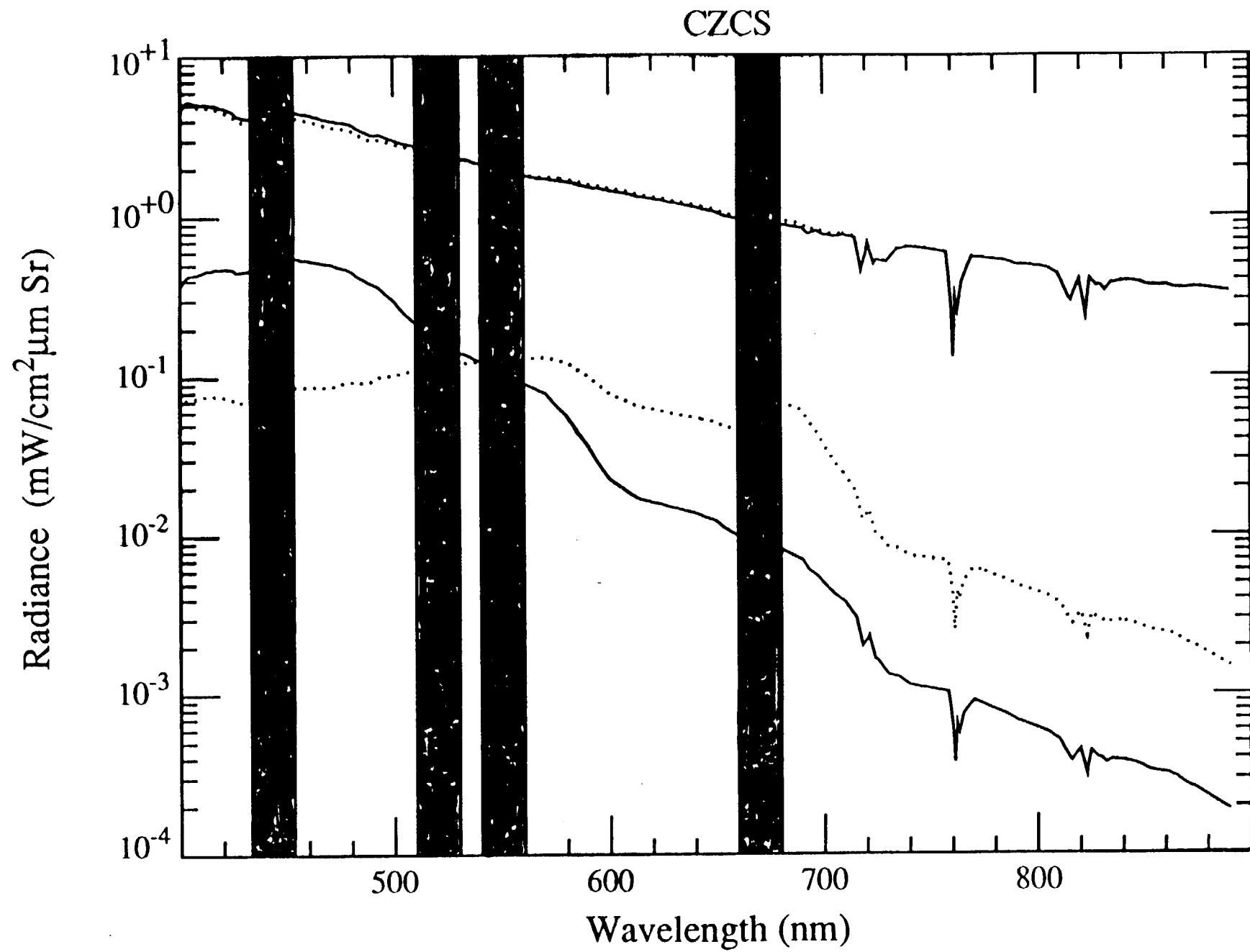
**University of Miami**

**APRIL 15, 1992**

## **OUTLINE**

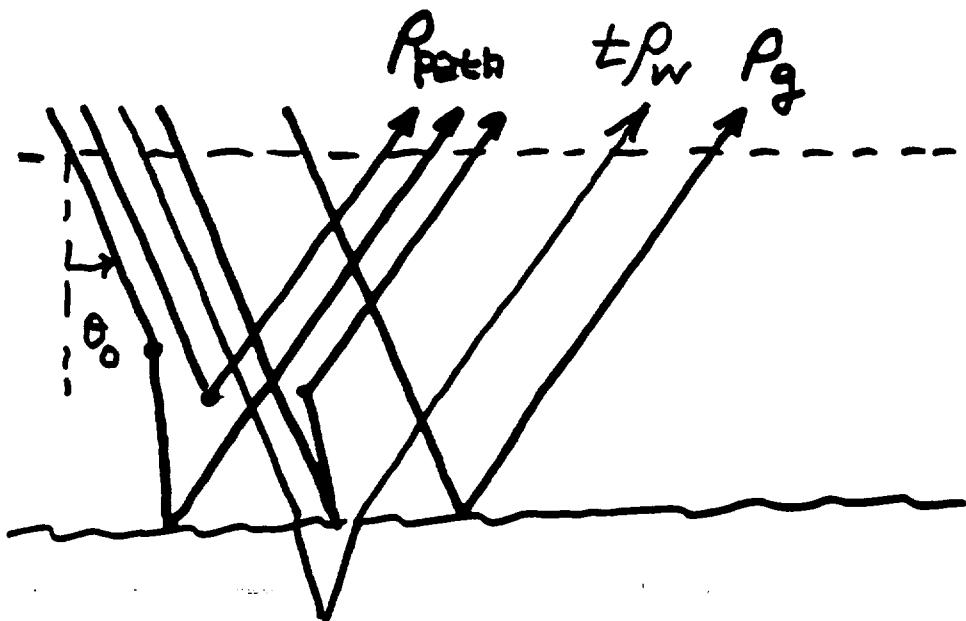
- 1. INTRODUCTION**
- 2. REVIEW OF FIRST-ORDER CORRECTION (CZCS)**
- 3. PROPOSED SECOND-ORDER ALGORITHM (SEAWIFS)**
- 4. IMPLEMENTATION**





## Atmospheric Correction

$$\rho_t = \rho_{\text{path}} + \rho_g + t \rho_w; \quad \rho = \frac{\pi L}{F_0 \cos \theta_0}$$

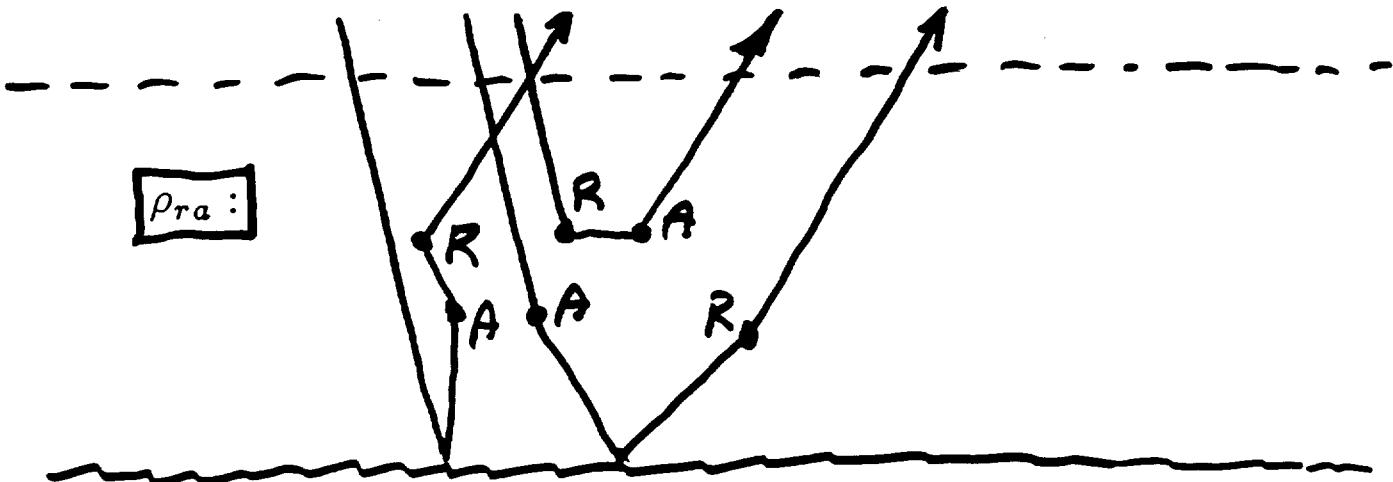
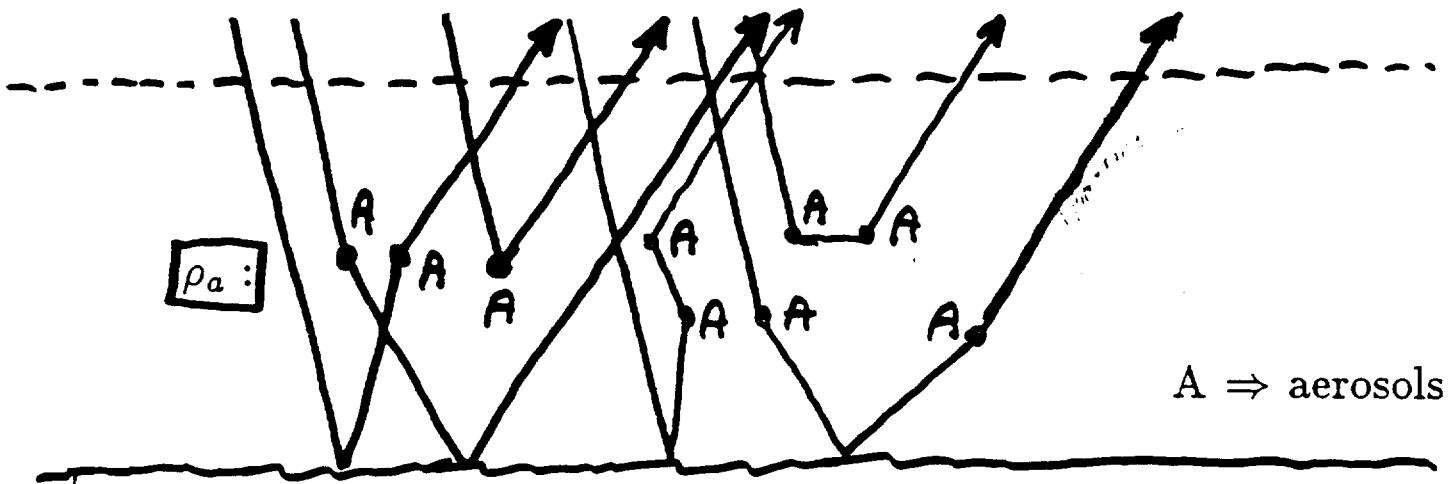
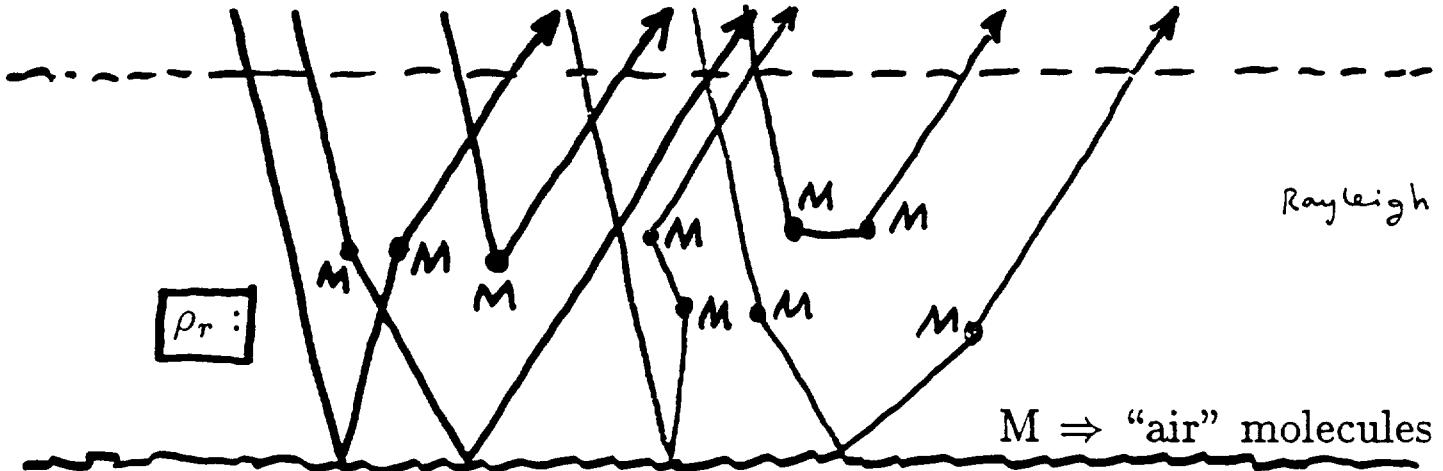


IGNORE  $\rho_g$ , THEN

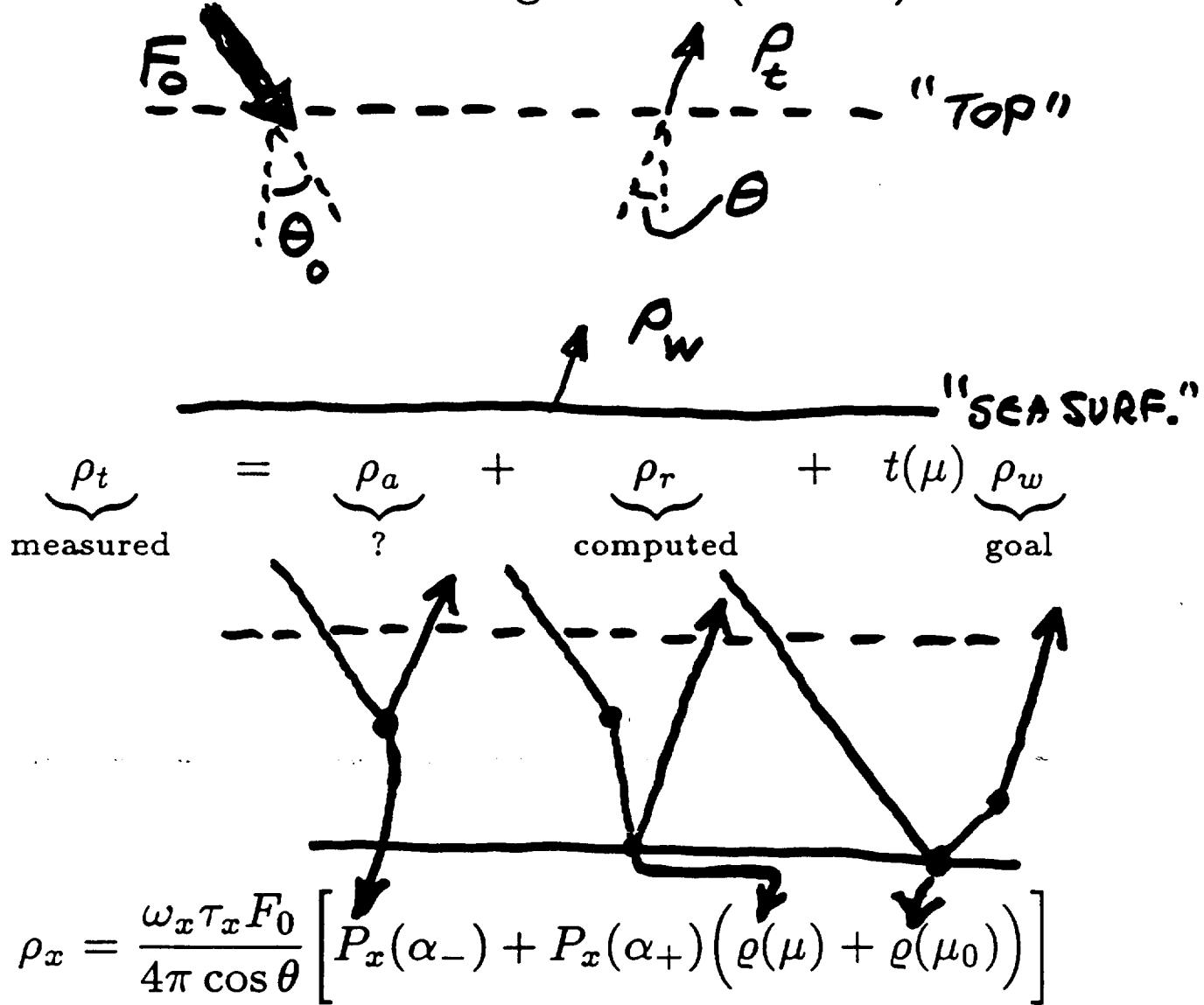
$$\rho_t = \rho_{\text{path}} + t \rho_w.$$

NOW LOOK AT THIS IN MORE DETAIL.

$$\rho_t = \rho_r + \rho_a + \rho_{ra} + t \rho_w.$$



## First-Order Algorithm (CZCS)



$$\cos \alpha_{\pm} = \pm \mu \mu_0 \sqrt{1 - \mu^2} \sqrt{1 - \mu_0^2} \cos \Delta \phi$$

$$x = r \quad \text{or} \quad a, \quad \mu = \cos \theta, \quad \mu_0 = \cos \theta_0$$

Define

$$S(\lambda, \lambda_0) = \frac{\rho_a(\lambda)}{\rho_a(\lambda_0)}$$

Then for single scattering

$$\begin{aligned} S &= \frac{\omega_a(\lambda)\tau_a(\lambda)F_0(\lambda)}{\omega_a(\lambda_0)\tau_a(\lambda_0)F_0(\lambda_0)} \left[ \frac{P_a(\lambda, \alpha_-) + P_a(\lambda, \alpha_+)[\varrho(\mu) + \varrho(\mu_0)]}{P_a(\lambda_0, \alpha_-) + P_a(\lambda_0, \alpha_+)[\varrho(\mu) + \varrho(\mu_0)]} \right] \\ &= \epsilon(\lambda, \lambda_0) \frac{F_0(\lambda)}{F_0(\lambda_0)} \end{aligned}$$

**NOTE:**

For a given aerosol “type”  $\epsilon(\lambda, \lambda_0)$  is independent of the aerosol **concentration**, and is almost independent of **position** over the image.

Consider two bands at  $\lambda$  and  $\lambda_0$ .  ~~$\rho_w(\lambda)$~~  = unknown

$$\begin{aligned}
 t(\lambda) \rho_w(\lambda) &= \rho_t(\lambda) - \rho_r(\lambda) - \underline{\rho_a(\lambda)} \\
 &= \rho_t(\lambda) - \rho_r(\lambda) - \underline{S(\lambda, \lambda_0)} \underline{\rho_a(\lambda_0)} \\
 &= \rho_t(\lambda) - \rho_r(\lambda) - \underline{S(\lambda, \lambda_0)} \left[ \rho_t(\lambda_0) - \rho_r(\lambda_0) - \underline{t(\lambda_0)} \underline{\rho_w(\lambda_0)} \right]
 \end{aligned}$$

where

$$t(\lambda) = \exp \left[ - \left( \tau_r/2 + \tau_{Oz} + (1 - \omega_a F) \tau_a \right) / \mu \right] \quad (< 1/6)$$

- (1) ignore  $(1 - \omega_a F) \tau_a$
- (2) choose  $\lambda_0$  such that  $\rho_w(\lambda_0) = 0$

$$\underline{\rho_w(\lambda)} = t(\lambda)^{-1} \left( \rho_t(\lambda) - \rho_r(\lambda) - \underline{S(\lambda, \lambda_0)} \left[ \rho_t(\lambda_0) - \rho_r(\lambda_0) \right] \right)$$


 $\epsilon(\lambda, \lambda_0) \frac{F_0(\lambda)}{F_0(\lambda_0)}$

## Apply to CZCS

**Goal:** Determine  $\rho_w$  at 443, 520, and 550 nm in order to estimate the pigment concentration.

**Problem:** Only *one* band for which  $\rho_w \approx 0$  and we need two.

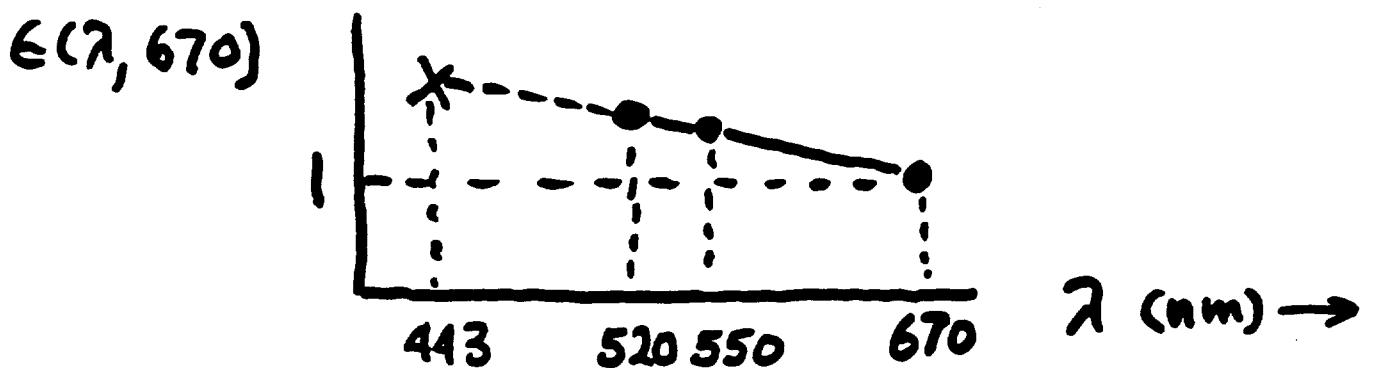
For low pigments can use Clear Water Radiance Concept:

When  $C \leq 0.25 \text{ mg/m}^3$ ,  $[\rho_w]_n$  in

$$\rho_w(\lambda) = [\rho_w(\lambda)]_N \cos \theta_0 \exp \left[ -(\tau_r/2 + \tau_{Oz}) / \cos \theta_0 \right],$$

is independent of  $C$  at 520, 550, and 670 nm.

Thus, use “clear water” regions to find  $\epsilon(520, 670)$ ,  $\epsilon(550, 670)$ , and  $\epsilon(670, 670)$ . Extrapolate these to find  $\epsilon(443, 670)$  and use this  $\epsilon$  set throughout the entire image.



Note: This hinges on  $\epsilon(\lambda, \lambda_0) \approx \text{const.}$  for given *aerosol type*.

SLIDES

TO SHOW  $\epsilon \approx \text{const}$

## **Difficulties:**

- 1) May be no “clear water” in the image of interest.
- 2) The aerosol type may vary over the image causing the  $\epsilon$ 's to vary.
- 3) The aerosol phase function depends weakly on wavelength which implies that the  $\epsilon$ 's will depend on position in the image even if all of the other approximations (single scattering) are valid.
- 4) It also ignores several processes:
  - a) Multiple Scattering
  - b) Surface Roughness
  - c) Vertical Structure
  - d) Whitecaps
  - e) Variations in Ozone
  - f) Variations in Pressure
  - g) Variations  $\rho_w$  with Viewing Angle

**CZCS ALGORITHM ERRORS**  
**(OVER AND ABOVE ERRORS IN THE  $\epsilon$ 'S)**

ASSUMPTION	$t(\lambda)\Delta\rho(\lambda)$	$\lambda$
$\rho_{ra} = 0$	$\sim 0.002$ ( $\tau_a \approx 0.3$ )	443
FLAT OCEAN	$\sim 0.0007$ (7.5 m/s)	443
$\Delta\tau_{oz} = 0$	$\sim 0.0008$ ( $\pm 50$ DU)	550
$\Delta P_0 = 0$	$\sim 0.002$ ( $\pm 15$ MB)	412
NO WHITECAPS	$\sim 0.002$ (7.5 m/s)	412–865

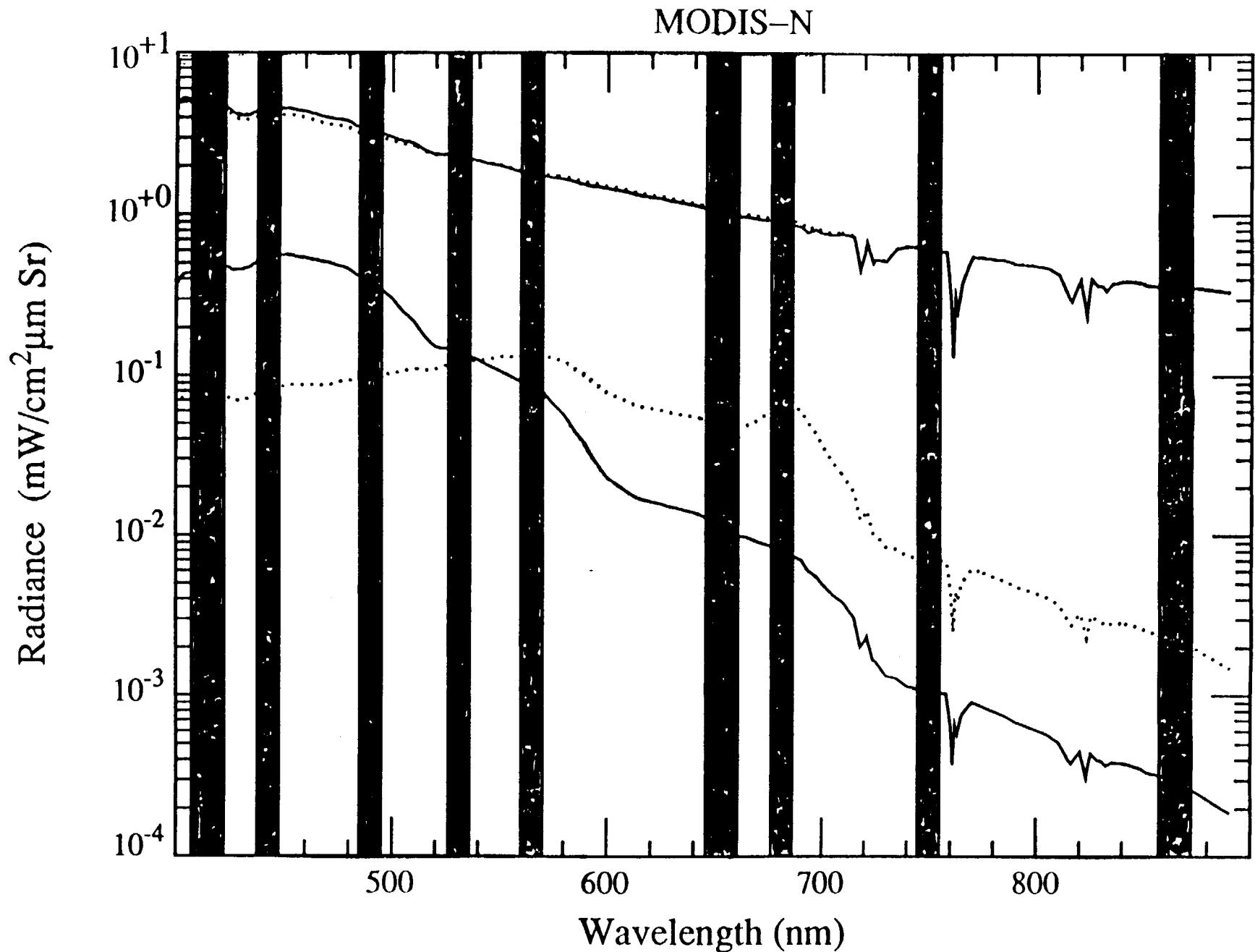
NOTE:

$$NE\Delta\rho_{CZCS}(443) = 0.0011$$

$$NE\Delta\rho_{CZCS}(550) = 0.00064$$

$$NE\Delta\rho_{SeaWiFS}(443) = 0.00049$$

$$NE\Delta\rho_{SeaWiFS}(550) = 0.00031$$



# SeaWiFS

## PERFORMANCE FOR $\theta_0 = 60^\circ$ AT THE SCAN EDGE

BAND	$\lambda$	$\rho_{\max}$	$\rho_{\text{typ}}$	$\rho_w$	$\text{NE}\Delta\rho$	CZCS
1	402–422	0.50	0.34	0.040	0.00068	—
2	433–453	0.46	0.29	0.038	0.00049	0.0011
3	480–500	0.36	0.23	0.024	0.00035	—
4	510–530	0.32	0.19	0.0096	0.00031	0.00058
5	555–575	0.26	0.16	0.0040	0.00031	0.00064
6	655–675	0.17	0.105	0.0004	0.00024	0.00051
7	745–785	0.15	0.081	—	0.00017	—
8	845–785	0.13	0.069	—	0.00015	—

TO UTILIZE THE FULL SENSITIVITY OF SEAWiFS,  
 THE ERROR IN ATMOSPHERIC CORRECTION SHOULD  
 BE  $\lesssim 0.0003 - 0.0007$ .

NOTE:

$$\text{NE}\Delta\rho_{\text{SeaWiFS}} \approx \frac{1}{2} \text{NE}\Delta\rho_{\text{CZCS}}$$

$$\text{NE}\Delta\rho_{\text{MODIS}} \approx \frac{1}{2} \text{NE}\Delta\rho_{\text{SeaWiFS}}$$

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## EXAMPLE OF ALGORITHM PERFORMANCE

EXAMINE THE ALGORITHM IN THE FOLLOWING MANNER:

1. ASSUME CZCS BANDS AT 443, 565, AND 665 NM.
2. ASSUME  $C < 0.25 \text{ MG/M}^3 \implies \rho_w(565) \text{ AND } \rho_w(665)$  AND KNOWN.
3. USE ALGORITHM TO FIND  $\rho_w(443) \implies C$ .

IN SUCH A SCENARIO, WE KNOW  $C < 0.25 \text{ MG/M}^3$  AND ARE TRYING TO FIND ITS ACTUAL VALUE. THIS SITUATION OFTEN OCCURS WITH CZCS, E.G., THE SARGASSO SEA IN SUMMER.

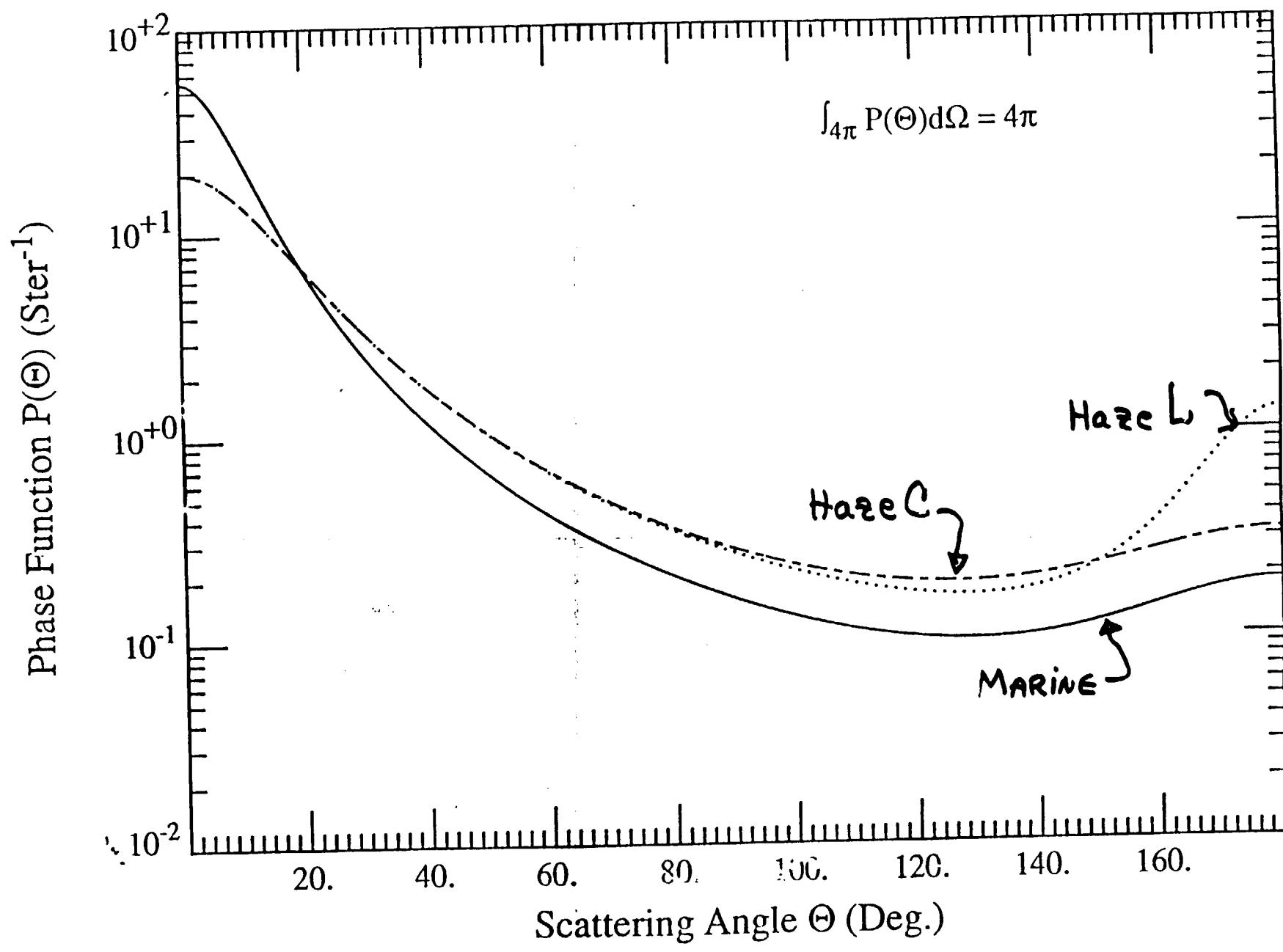
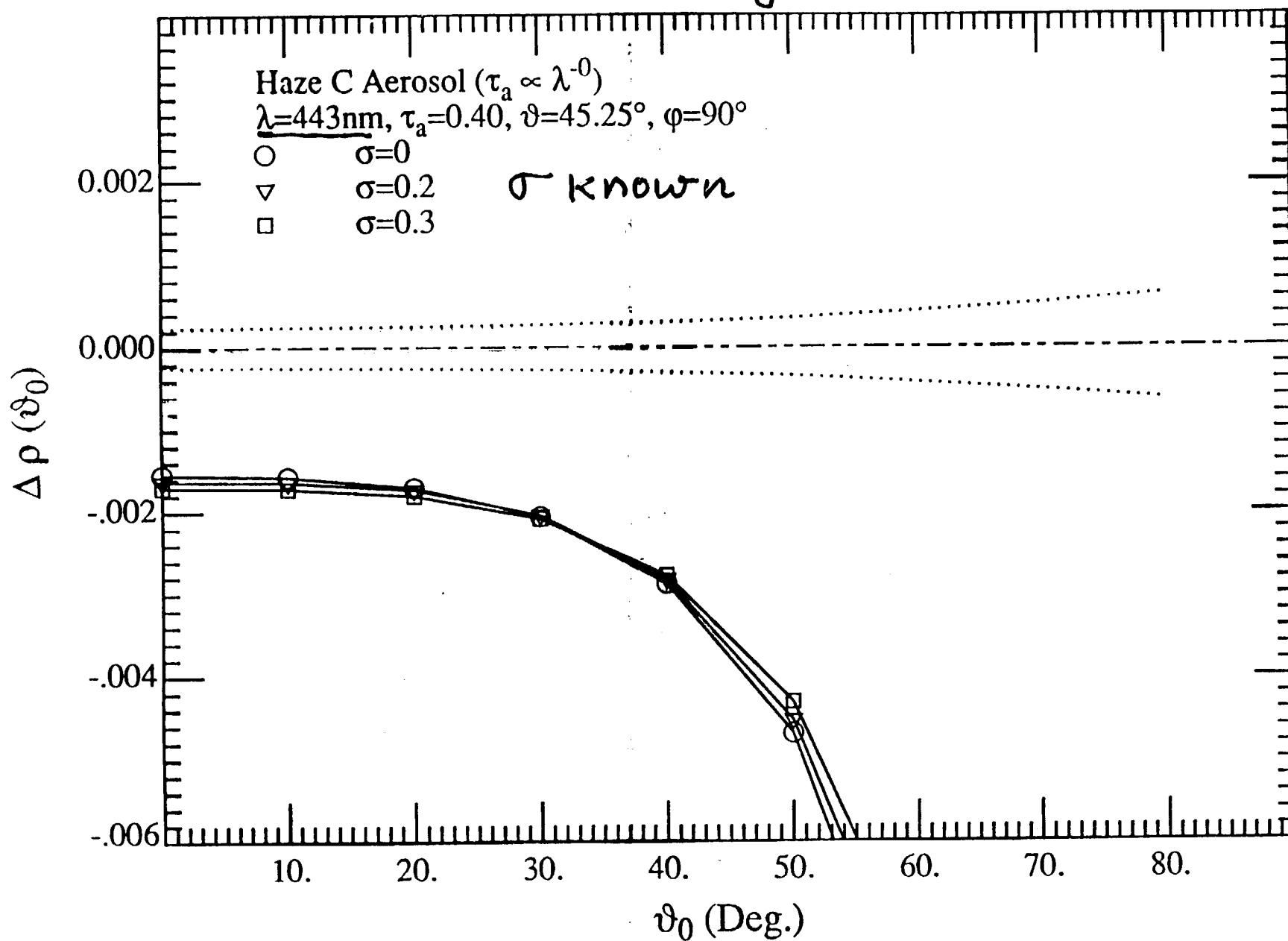
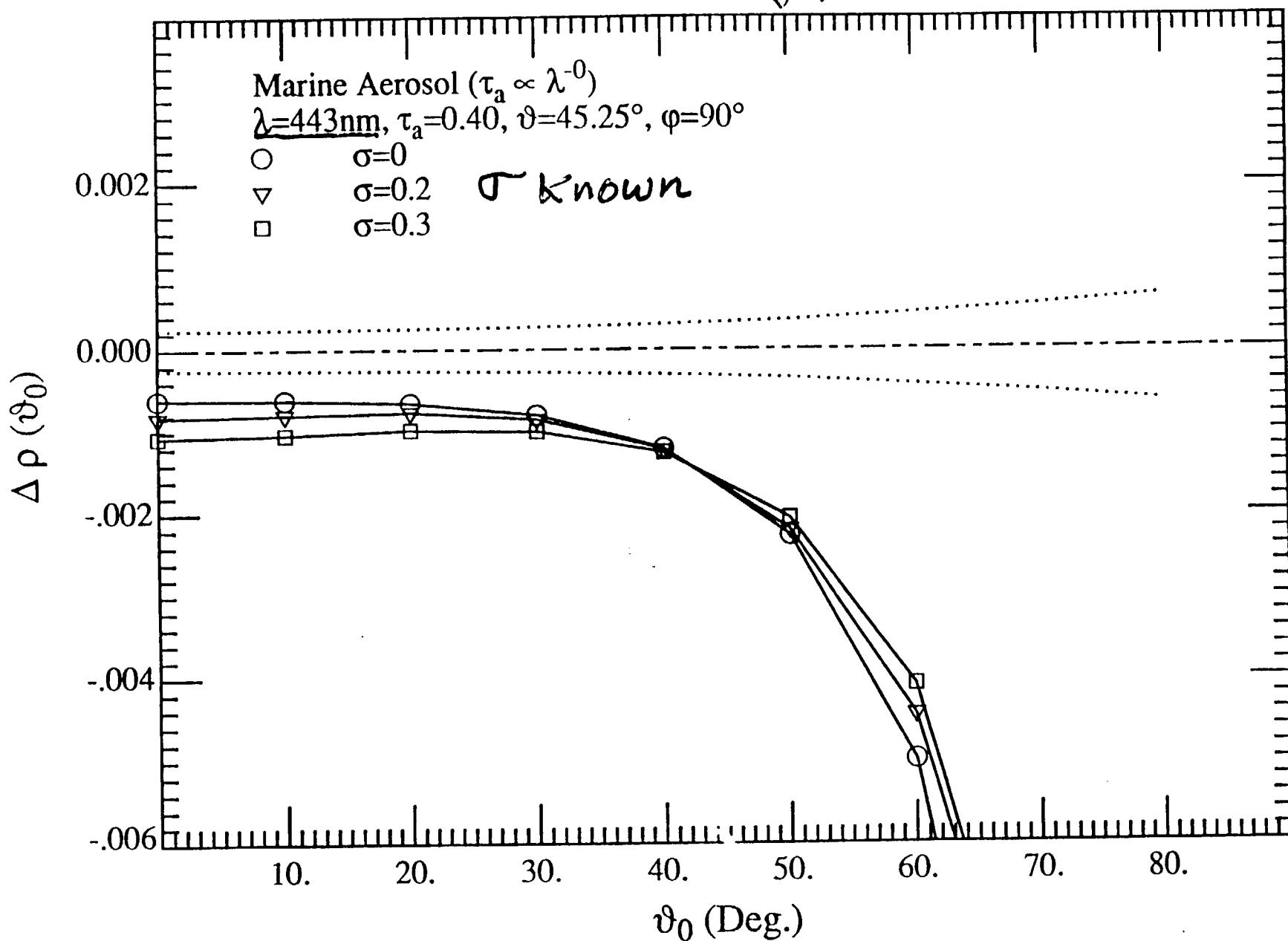


Figure 2

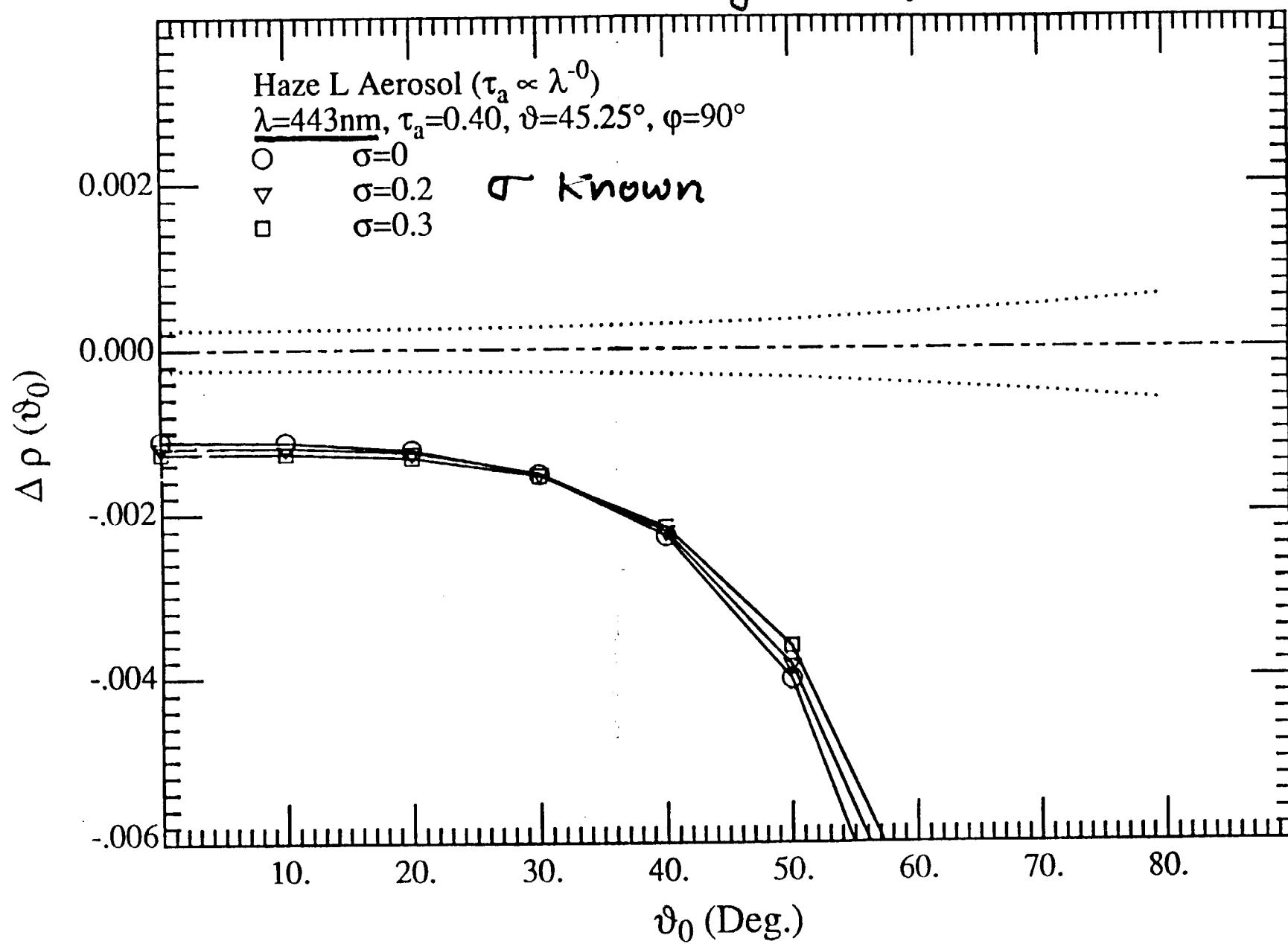
# CZCS Algorithm



# CZCS Algorithm



# CZCS Algorithm



## Sketch of Preliminary Algorithm Ideas

ASSUME  $\rho_w$  IS KNOWN AT 565 AND 665 AND DESIRED AT 443.

COMPUTE  $\rho_t$  INCLUDING ALL PROCESSES AND  $\rho_r$  WHICH IS  $\rho_t$  WHEN THERE IS NO AEROSOL.

THEN SINCE

$$\rho_t = \rho_r + \rho_a + \rho_{ra} + t\rho_w$$

WE HAVE

$$\rho_a + \rho_{ra} = \rho_t - \rho_r - t\rho_w$$

WHICH PROVIDES  $\rho_a + \rho_{ra}$  AT 565 AND 665 NM.

FROM THESE WE MUST ESTIMATE  $\rho_a + \rho_{ra}$  AT 443 NM.

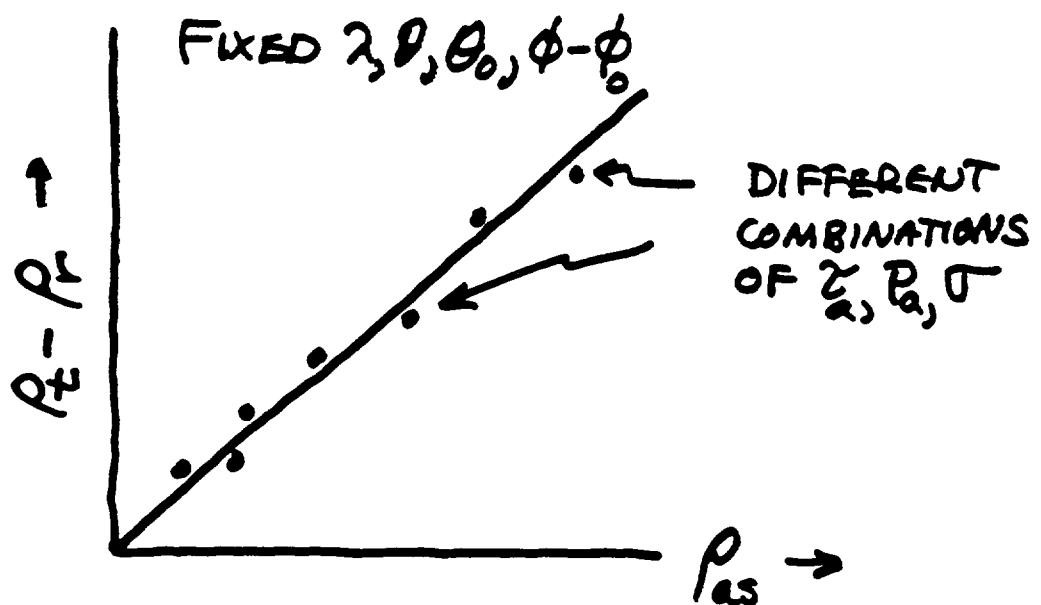
TO DO THIS WE FOLLOW GORDON AND CASTANO (1989) AND RELATE  $\rho_a + \rho_{ra}$  TO  $\rho_{as}$  — THE AEROSOL REFLECTANCE IN THE SINGLE SCATTERING/FLAT OCEAN APPROXIMATION:

$$\rho_{as} = \omega_a(\lambda)\tau_a(\lambda)p_a(\theta, \theta_0, \lambda)/4\pi \cos \theta \cos \theta_0,$$

WHERE

$$p_a(\theta, \theta_0, \lambda) = P_a(\theta_-, \lambda) + (\rho(\theta) + \rho(\theta_0))P_a(\theta_+, \lambda),$$

$$\cos \theta_{\pm} = \pm \cos \theta_0 \cos \theta - \sin \theta_0 \sin \theta \cos(\phi - \phi_0).$$

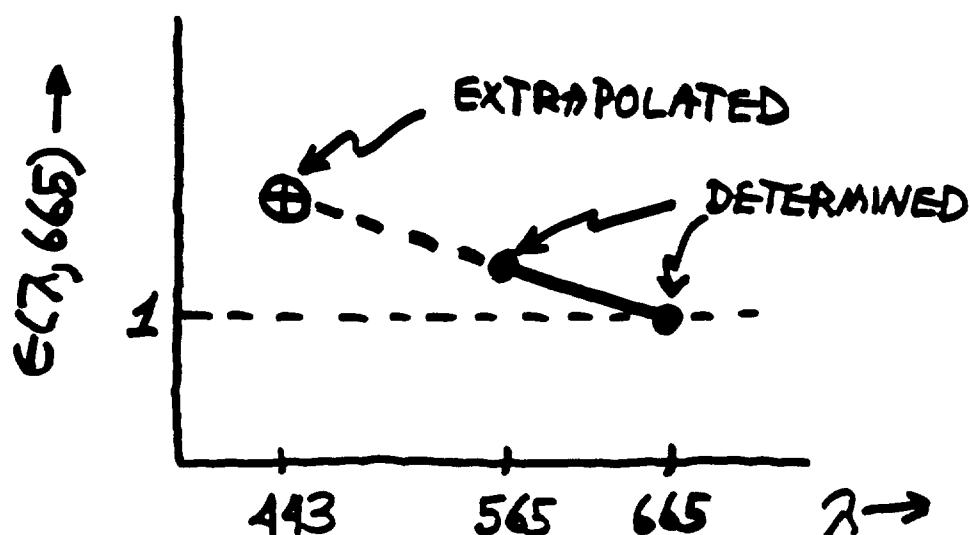


$$\rho_t(\lambda) - \rho_r(\lambda) = C_1(\lambda) + C_2(\lambda)\rho_{as}(\lambda)$$

THEN,

$$\frac{\rho_{as}(\lambda_i)}{\rho_{as}(\lambda_j)} = \frac{\omega_a(\lambda_i)\tau_a(\lambda_i)p_a(\theta, \theta_0, \lambda_i)}{\omega_a(\lambda_j)\tau_a(\lambda_j)p_a(\theta, \theta_0, \lambda_j)} \equiv \epsilon(\lambda_i, \lambda_j)$$

NOW, SINCE  $\rho_w(565)$  AND  $\rho_w(665)$  ARE KNOWN,  $\rho_{as}(565)$  AND  $\rho_{as}(665)$  CAN BE DETERMINED. THESE PROVIDE  $\epsilon(565, 665)$  AND  $\epsilon(665, 665)$ , AND WE EXTRAPOLATE TO FIND  $\epsilon(443, 665)$



**FINALLY,**

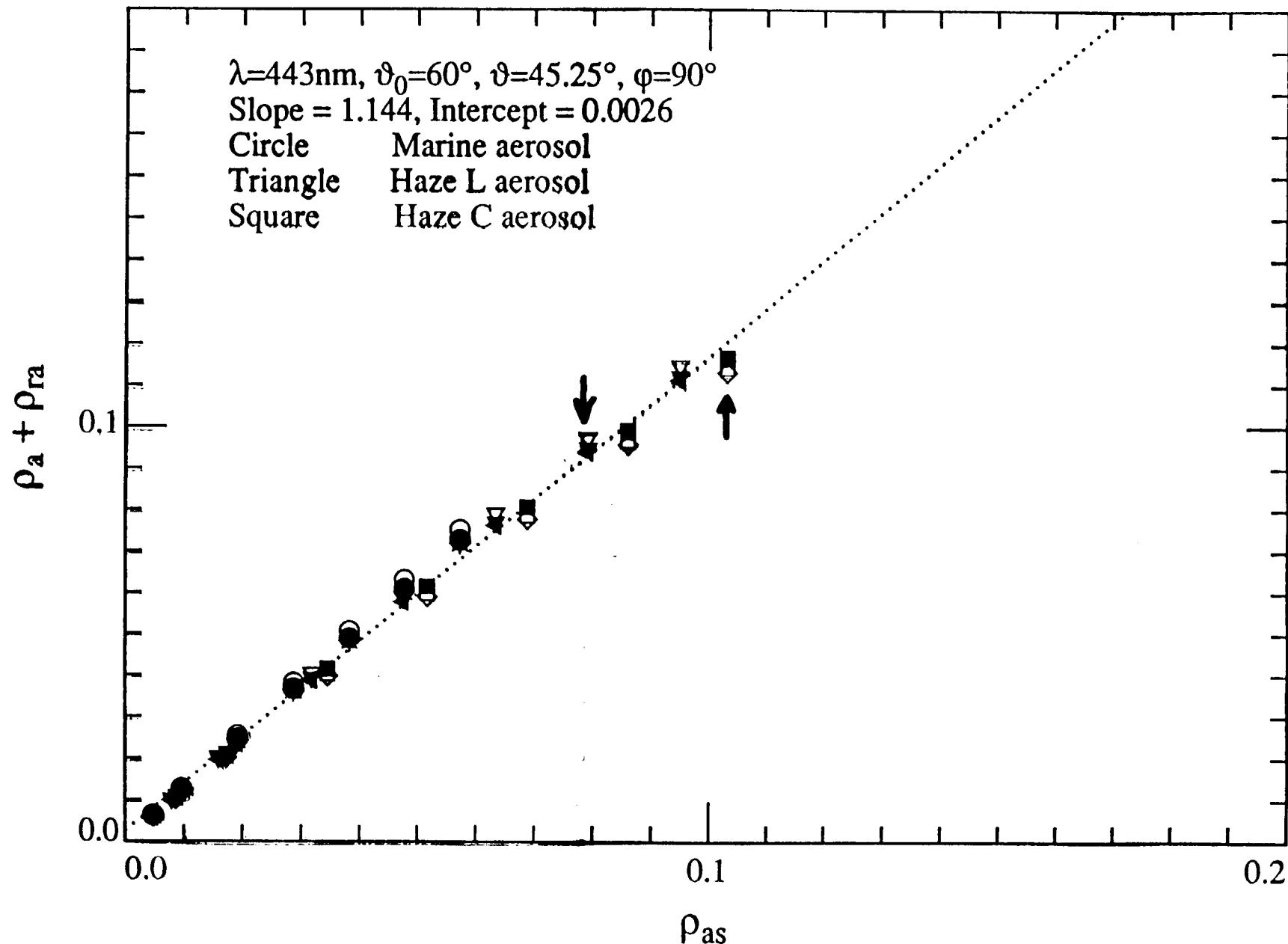
$$\rho_{as}(443) = \epsilon(443, 665)\rho_{as}(665)$$

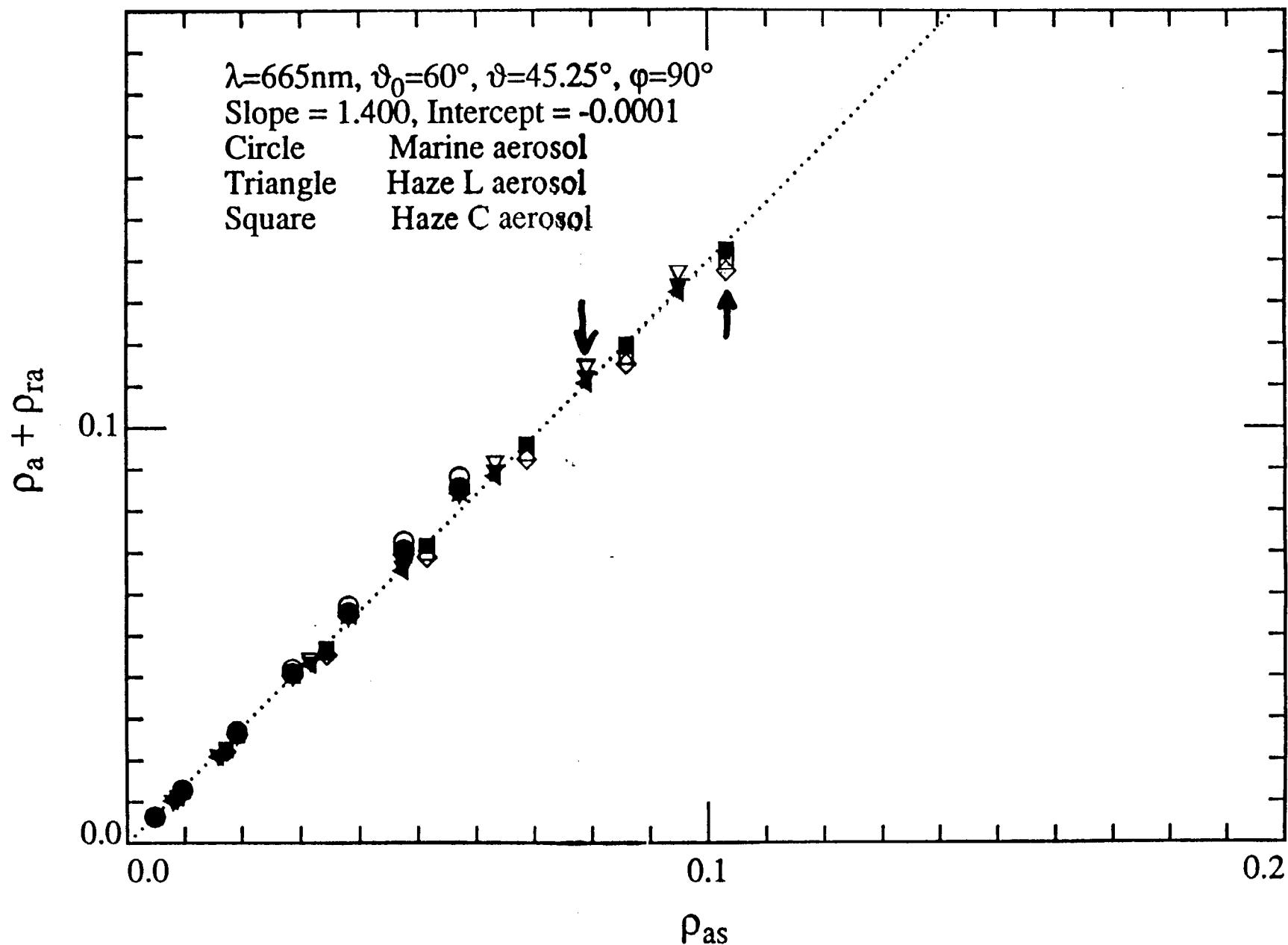
**AND**

$$\rho_a(443) + \rho_{ra}(443) = C_1(443) + C_2(443) \rho_{as}(443)$$

$t(443)\rho_w(443) = \rho_t(443) - \rho_r(443) - \rho_a(443) - \rho_{ra}(443)$

The diagram illustrates the decomposition of total density into measured and calculated components. At the top, the equation  $\rho_a(443) + \rho_{ra}(443) = C_1(443) + C_2(443) \rho_{as}(443)$  is shown. A bracket underlines the terms  $\rho_a(443)$  and  $\rho_{ra}(443)$ , and another bracket underlines the terms  $C_1(443)$  and  $C_2(443) \rho_{as}(443)$ . A curved arrow points from the bracketed terms down to the corresponding terms in the equation below. The equation below is  $t(443)\rho_w(443) = \rho_t(443) - \rho_r(443) - \rho_a(443) - \rho_{ra}(443)$ . Below this equation, two arrows point upwards from the words "MEASURED" and "CALCULATED" to the terms  $\rho_t(443) - \rho_r(443)$  and  $\rho_a(443) - \rho_{ra}(443)$  respectively.





## TEST OF PROPOSED ALGORITHM

TEST ALGORITHM VIA SIMULATION IN THE SAME MANNER AS BEFORE:

1. ASSUME THE SENSOR IS A MORE SENSITIVE CZCS, I.E., BANDS AT 443, 565, AND 665 NM.
2. ASSUME  $C < 0.25 \text{ MG/M}^3 \implies \rho_w(565) \text{ AND } \rho_w(665) \text{ AND KNOWN.}$
3. USE ALGORITHM TO FIND  $\rho_w(443) \implies C.$

IN SUCH A SCENARIO, WE KNOW  $C < 0.25 \text{ MG/M}^3$  AND ARE TRYING TO FIND ITS ACTUAL VALUE. THIS SITUATION OFTEN OCCURS WITH CZCS, E.G., THE SARGASSO SEA IN SUMMER.

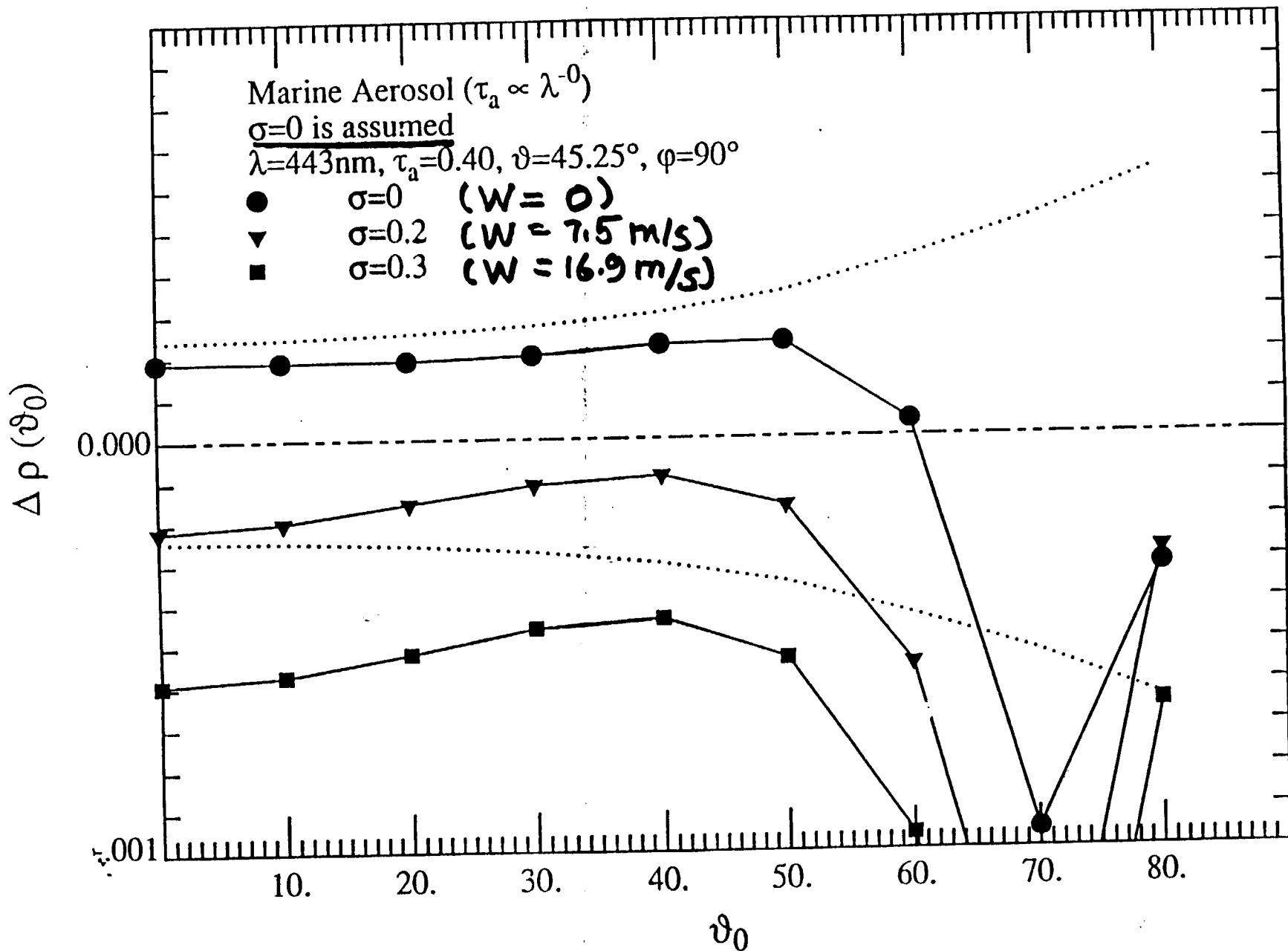


Figure 6.

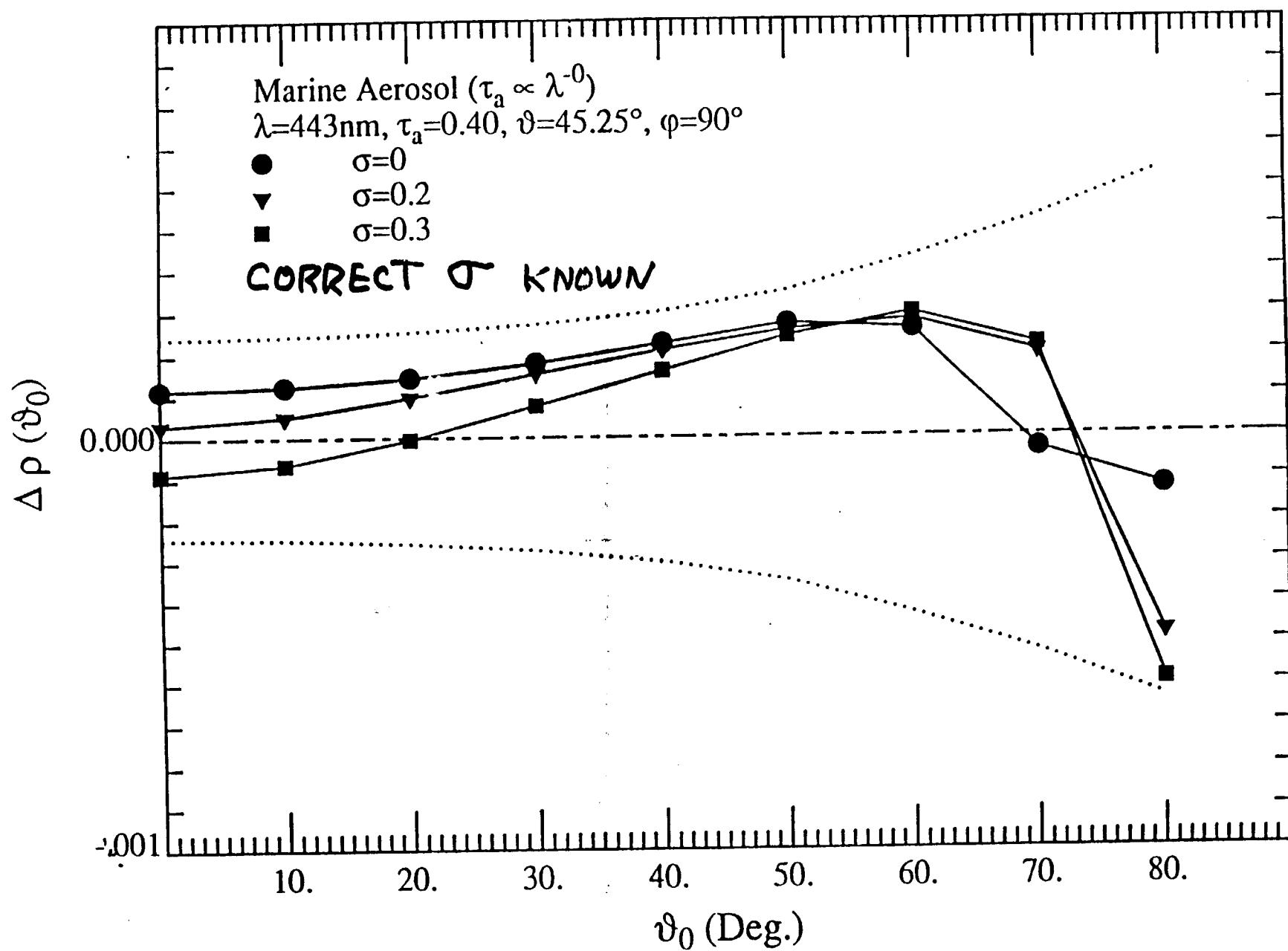


Figure 7

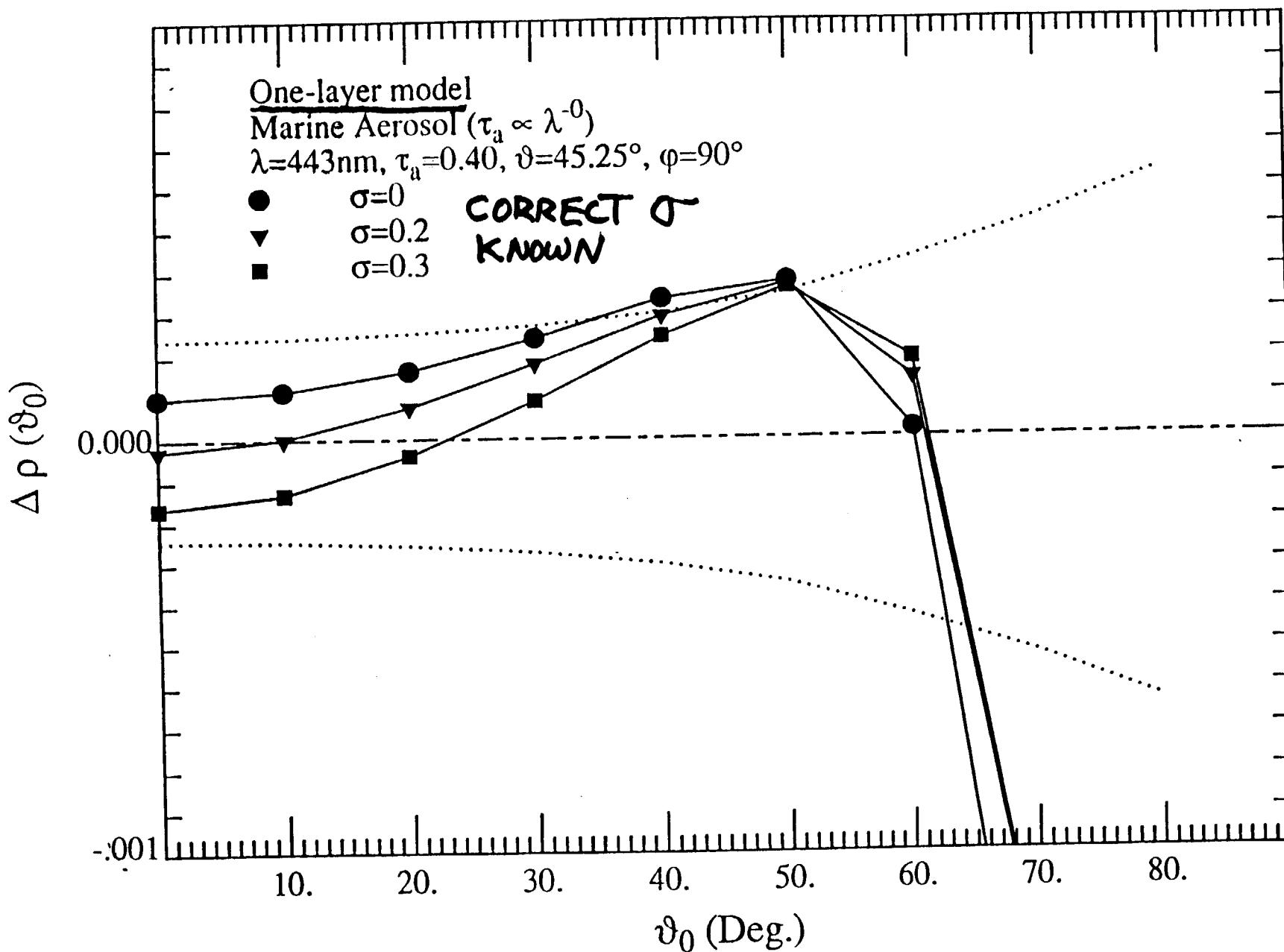


Fig. 2 10.

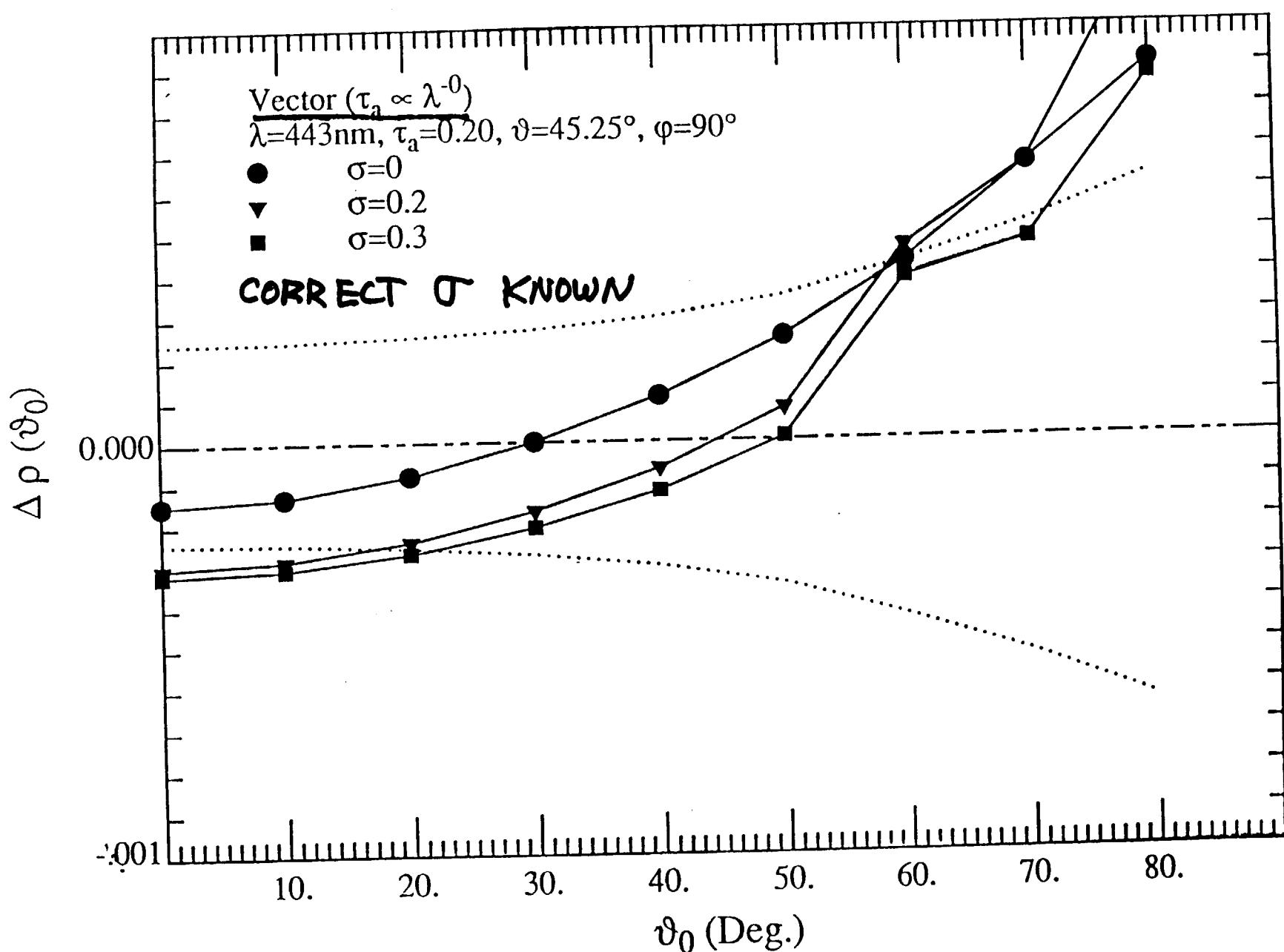
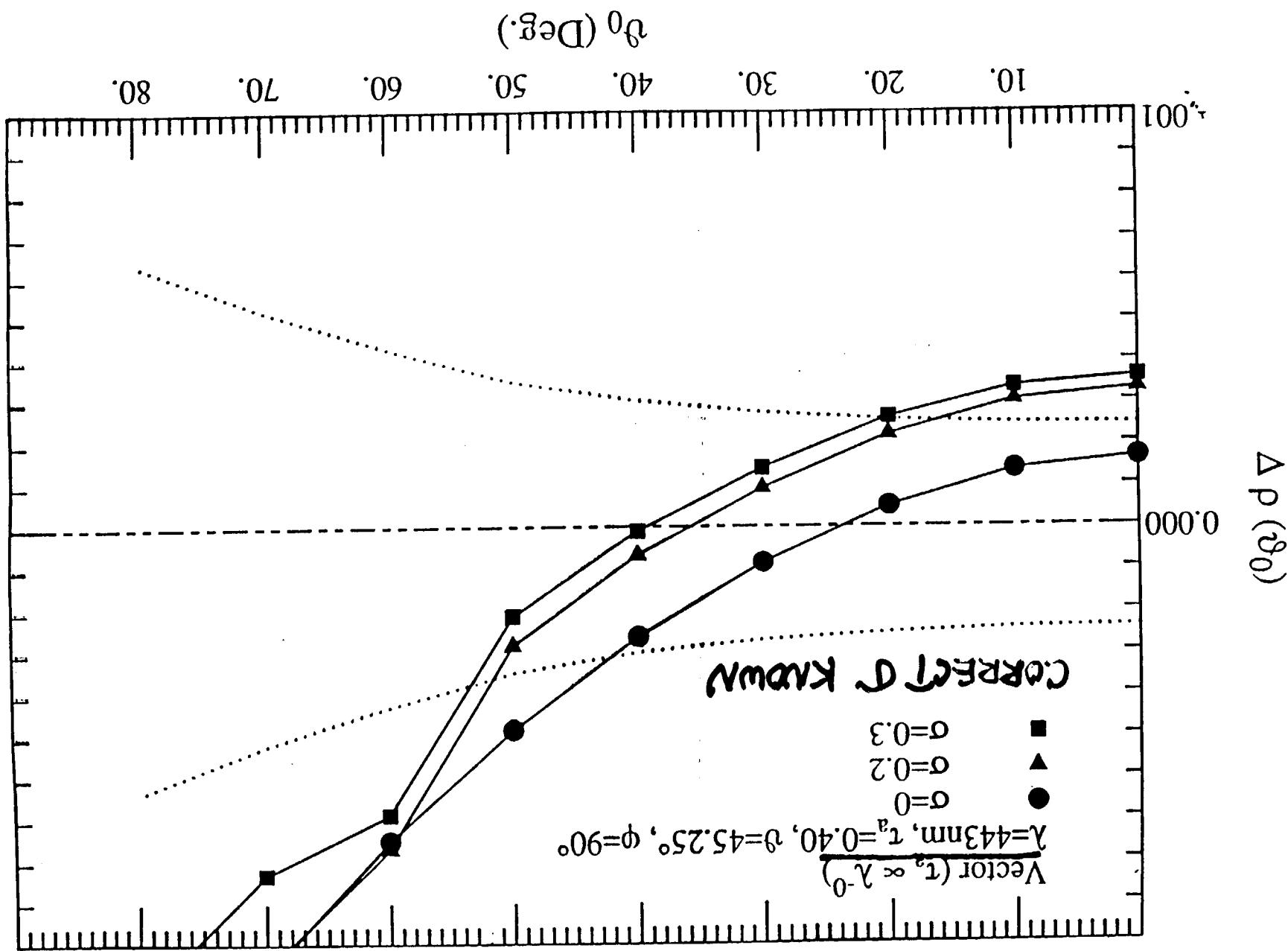
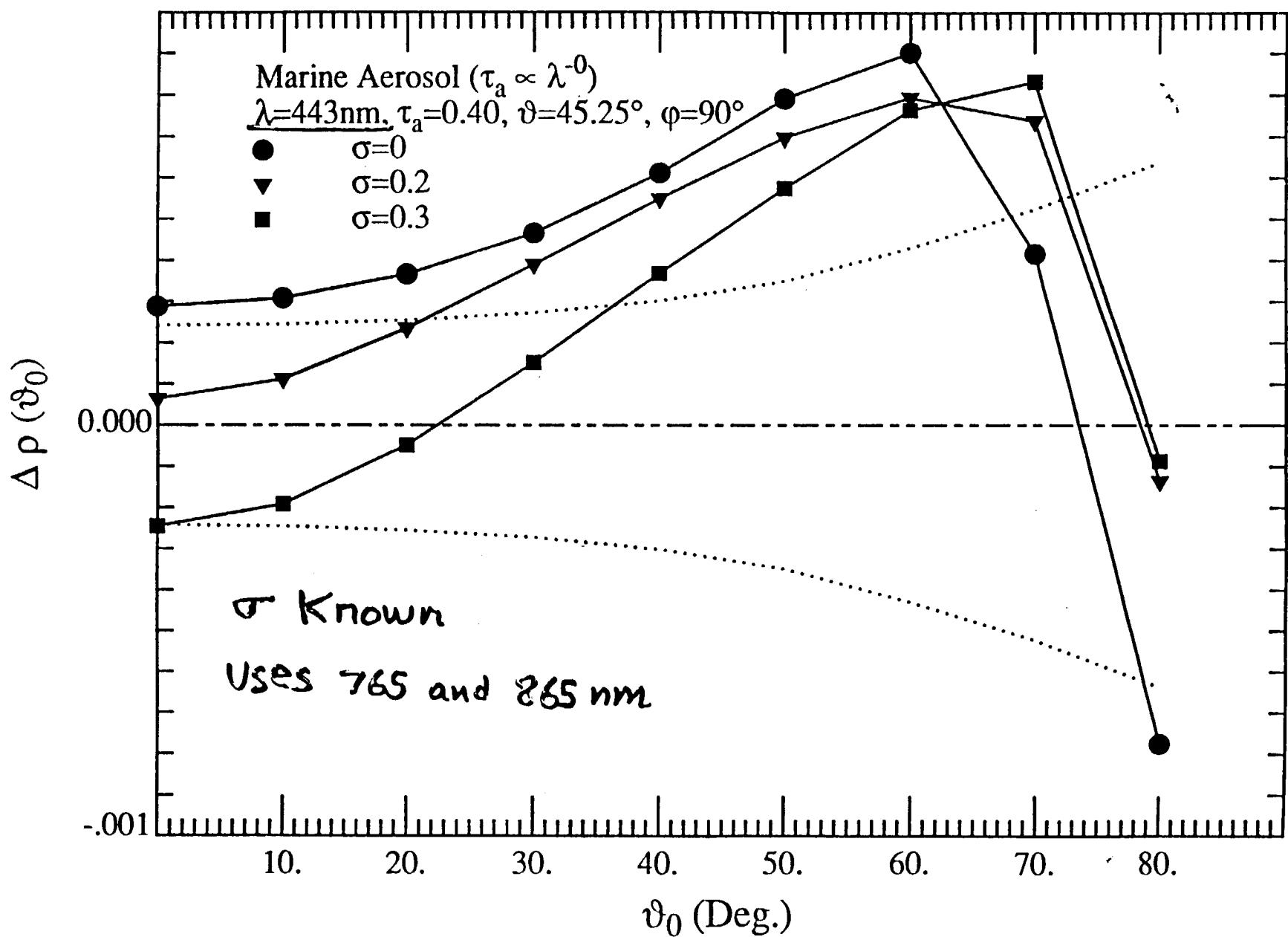
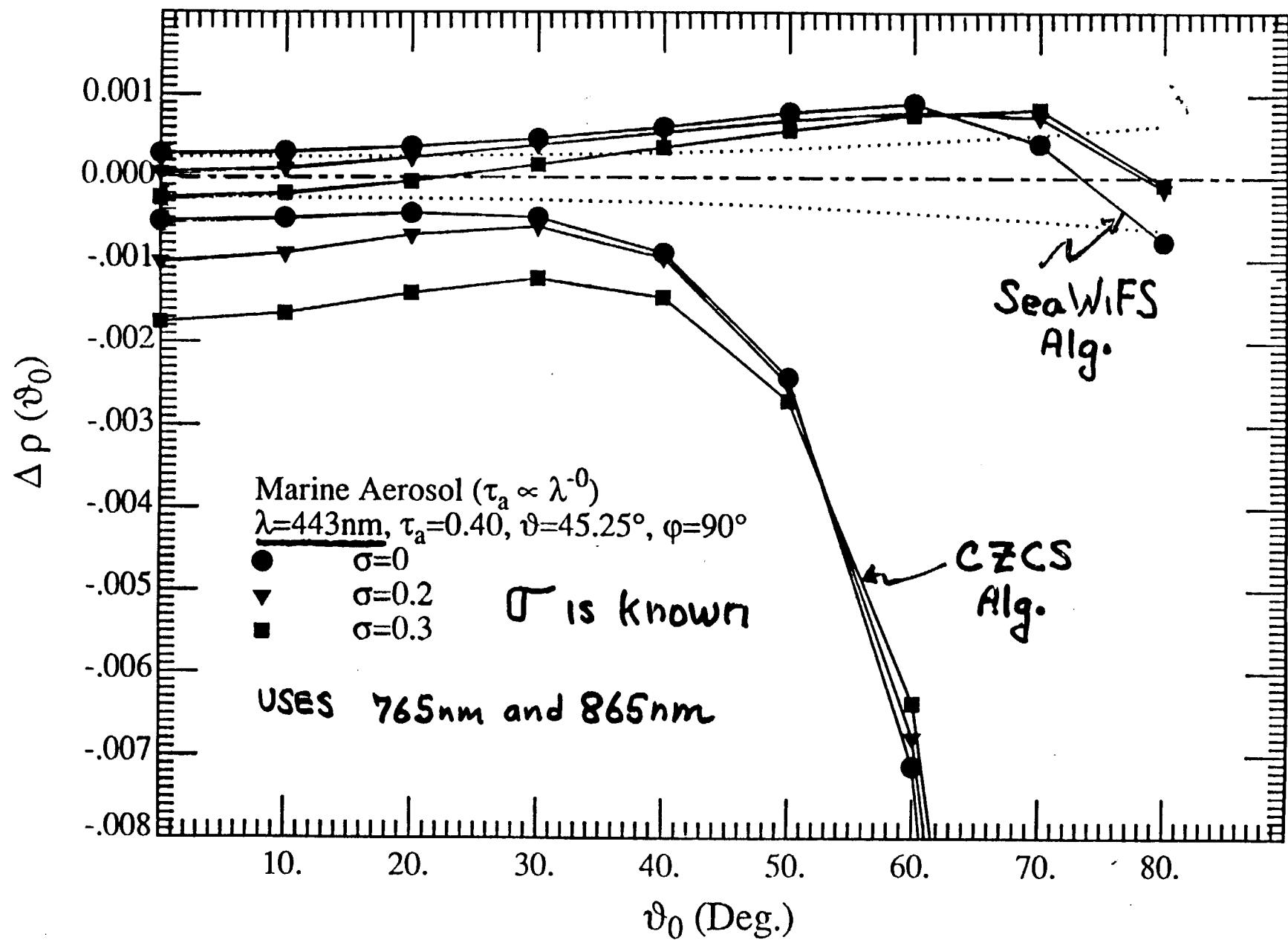


Figure 17

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## WHAT'S NEXT ?

1. DEVELOP CODE FOR TESTING ON CZCS DATA
  - A. LOW LATITUDE — EXPECT LITTLE DIFFERENCE
  - B. HIGH LATITUDE — EXPECT IMPROVEMENT
  - C. CZCS CALIBRATION COULD BE A PROBLEM
2. IMPLEMENTATION FOR SEAWIFS PROCESSING
  - A. NEED  $\tau_{Oz}$ ,  $P_0$ , WINDS (R. EVANS)
  - B. NEED  $\rho_r$  FOR VARIOUS WIND SPEEDS
  - C. NEED  $C_1$  AND  $C_2$  FOR ALL SCENARIOS
  - D. NEED TO STUDY  $\epsilon$  EXTRAPOLATION
3. IS IT VIABLE FOR MODIS?
  - A. INVESTIGATE MORE CASES
  - B. COMPUTE  $C_1$  AND  $C_2$  USING VECTOR THEORY
  - C. SEAWIFS EXPERIENCE
4. WHAT MORE IS REQUIRED FOR MODIS?
  - A. DEVELOP AN ACCURATE WHITECAP MODEL
  - B. LOOK AT VARIATION OF  $\rho_w$  WITH VIEW ANGLE
  - C. EARTH CURVATURE AFFECTS
  - D. REMOVE RESIDUAL INSTRUMENT POLARIZATION